

# 26: Connecting the Electric Field and the Potential

Remember what the **Electric Field** and **Potential** are:

**They both describe how source charges communicate their electric effect over all space – see the PhET image**

$\vec{E}$  associates a **vector** with each point in space

$V$  associates a **scalar** with each point in space

They both describe the same thing, but in a slightly different language.

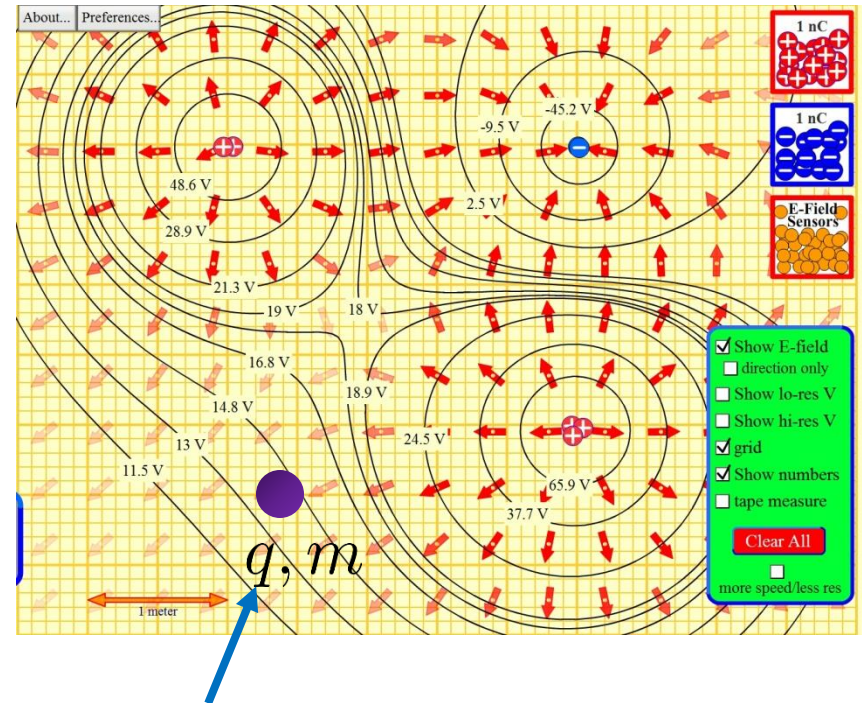
**What's the difference?**

Suppose we released a particle with **charge  $q$  and mass  $m$**  at some point. **How would we find its subsequent motion?**

**Using the Field:**  $\vec{F} = q\vec{E} \Rightarrow$  find the motion with Newton's 2<sup>nd</sup> Law,  $\vec{F} = m\vec{a}$

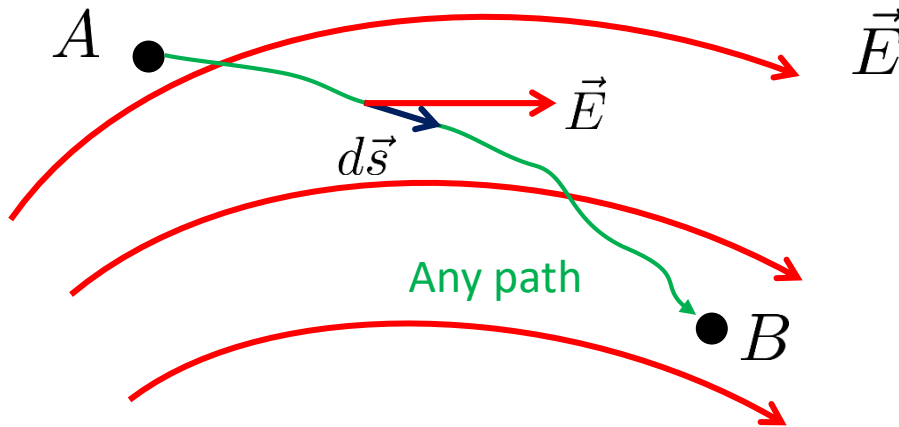
**Using the Potential:**  $U = qV \Rightarrow$  find the motion with Energy Conservation,  $\Delta E = \Delta K + \Delta U$

**How can we go back and forth between these two descriptions?**



# Going from the Electric Field to the Potential

*i.e.* If we know  $\vec{E}$  in a region of space, how do we find  $V$ ?



**What is the potential difference between points A and B?**

Recall from PHY181, the **Potential Energy** associated with a **Conservative Force** is:

$$\Delta U = - \int \vec{F} \cdot d\vec{s} \quad (\text{i.e. } \underline{\text{minus}} \text{ the work done by } F)$$

And for the **Electric Force**:  $\vec{F} = q\vec{E}$  and  $\Delta U = q\Delta V$

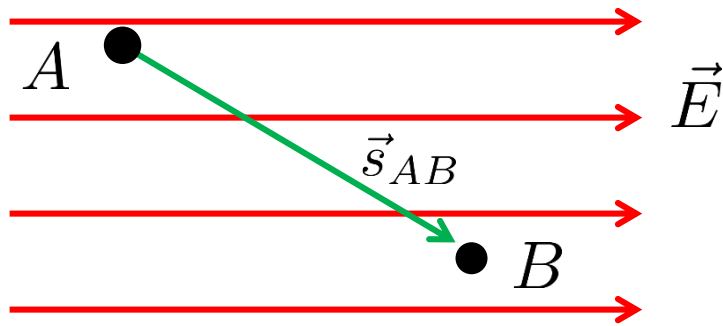
**So: the potential difference between two points is:**

$$\Delta V = V_B - V_A = - \int_A^B \vec{E} \cdot d\vec{s}$$

**Where the integral can be done over any path between A and B, since the electric force is conservative.**

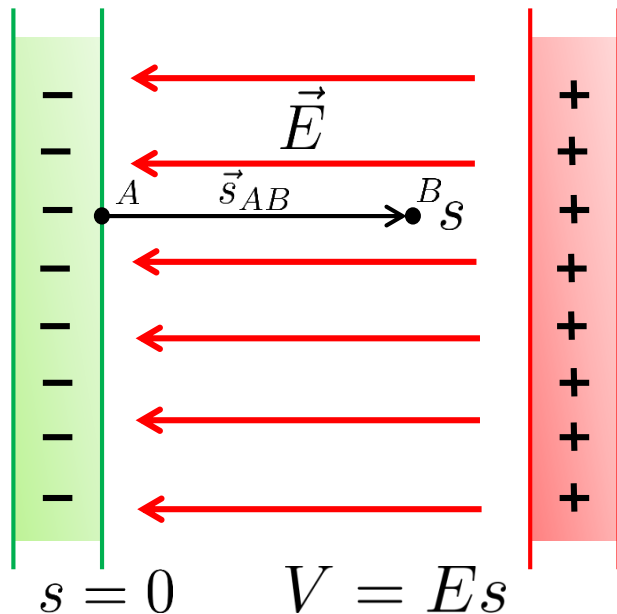
# Going from the Electric Field to the Potential

## For a Uniform Field



$$\begin{aligned}\Delta V &= V_B - V_A = - \int_A^B \vec{E} \cdot d\vec{s} \\ &= -\vec{E} \cdot \int_A^B d\vec{s} = -\vec{E} \cdot \vec{s}_{AB}\end{aligned}$$

Is this consistent with what we did last class?



$$\begin{aligned}\Delta V &= V_B - V_A = -\vec{E} \cdot \vec{s}_{AB} \\ &= -Es \cos 180^\circ \\ &= Es\end{aligned}$$

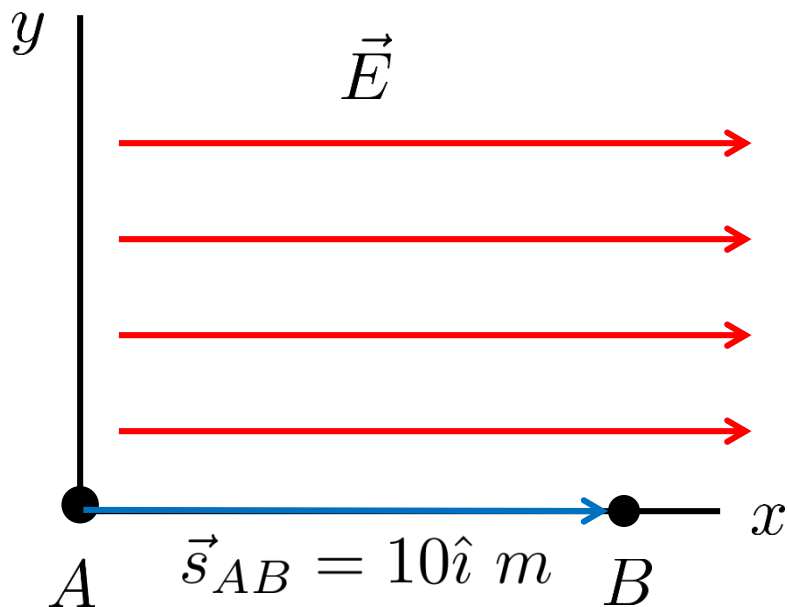
Perhaps you learned in a previous course that the **Potential is the field times the distance**, i.e.  $V = Ed$ . This is a special case. **The following whiteboard problems show cases where  $V = Ed$  isn't always right.**

# Whiteboard Problem 26-1: Potential from Field 1 (LC)

Find the potential difference,  $\Delta V = V_B - V_A = - \int_A^B \vec{E} \cdot d\vec{s}$ , between the origin,  $(x = 0, y = 0) \text{ m}$ , and point B at  $(x = 10, y = 0) \text{ m}$  for the electric field:

$$\vec{E} = 750\hat{i} \text{ V/m}$$

**First: Draw it:**



$$\begin{aligned}\Delta V &= V_B - V_A = - \int_A^B \vec{E} \cdot d\vec{s} \\ &= -\vec{E} \cdot \int_A^B d\vec{s} = -\vec{E} \cdot \vec{s}_{AB} \\ &= -(750 \text{ V/m})(10 \text{ m}) \\ &= -7500 \text{ V}\end{aligned}$$

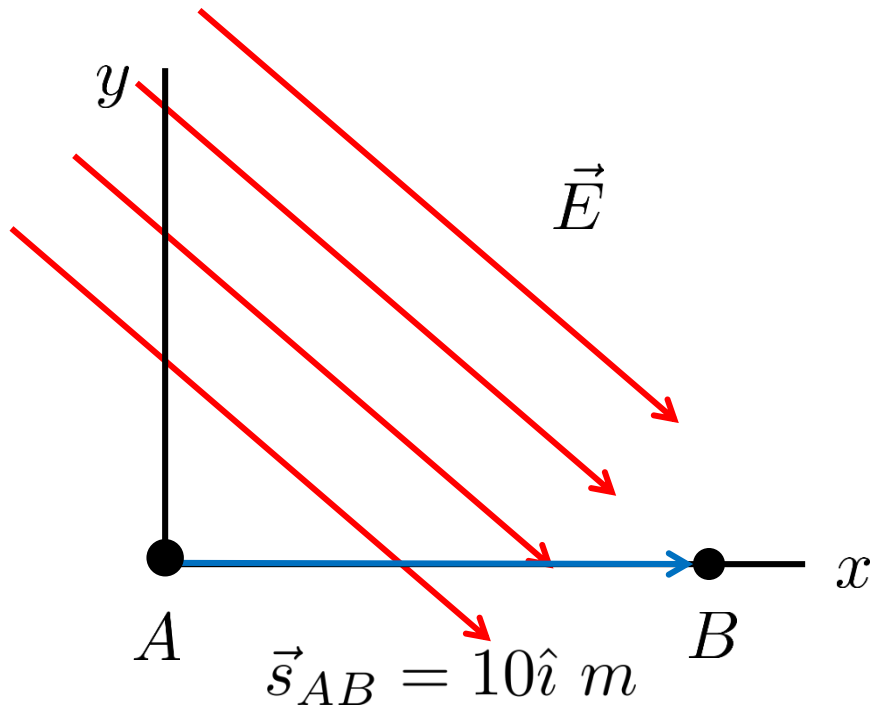
Note:  $V = Ed$  gives  $+7500 \text{ V}$ ; i.e. it doesn't get the sign right. For this field,  $V_B < V_A$ .

## Whiteboard Problem 26-1: Potential from Field 2 (LC)

Find the potential difference,  $\Delta V = V_B - V_A = -\int_A^B \vec{E} \cdot d\vec{s}$ , between the origin,  $(x = 0, y = 0) \text{ m}$ , and point B at  $(x = 10, y = 0) \text{ m}$  for the electric field:

$$\vec{E} = 750\hat{i} - 600\hat{j} \text{ V/m}$$

**First: Draw it:**



$$\begin{aligned}\Delta V &= V_B - V_A = -\int_A^B \vec{E} \cdot d\vec{s} \\ &= -\vec{E} \cdot \int_A^B d\vec{s} = -\vec{E} \cdot \vec{s}_{AB} \\ &= -(750 \text{ V/m})(10 \text{ m}) \\ &= -7500 \text{ V}\end{aligned}$$

(i.e. only that part of the displacement parallel to the field matters, the same as work done by a force.)

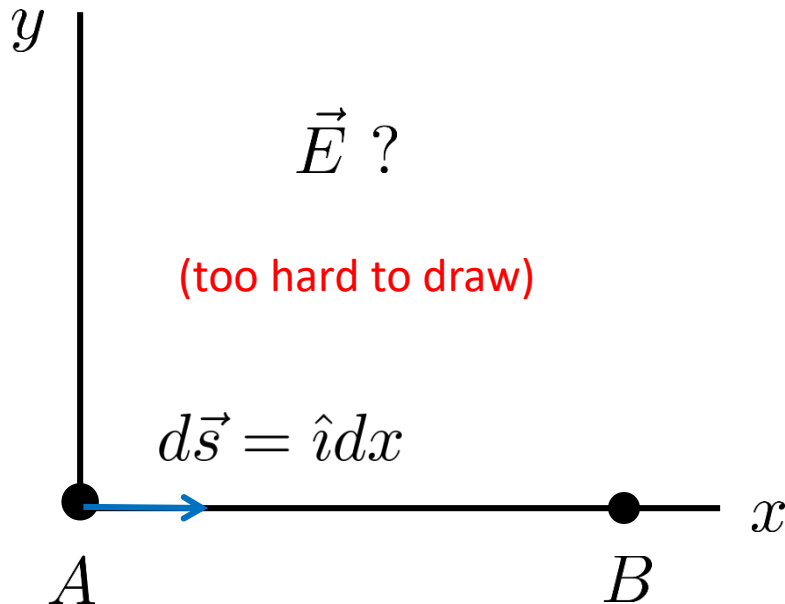
Note:  $V = Ed$  gives +9600 V.

## Whiteboard Problem 26-1: Potential from Field 3 (LC)

Find the potential difference,  $\Delta V = V_B - V_A = - \int_A^B \vec{E} \cdot d\vec{s}$ , between the origin,  $(x = 0, y = 0)$  m, and point B at  $(x = 10, y = 0)$  m for the electric field:

$$\vec{E} = 750x\hat{i} - 600y\hat{j} \text{ V/m}$$

First: Draw it:



$$\Delta V = V_B - V_A = - \int_A^B \vec{E} \cdot d\vec{s}$$

$$\vec{E} \cdot d\vec{s} = (750x)(dx)$$

$$\text{So, } \Delta V = - \int_0^{10} 750x dx$$

$$= -750 \left[ \frac{x^2}{2} \right]_0^{10}$$

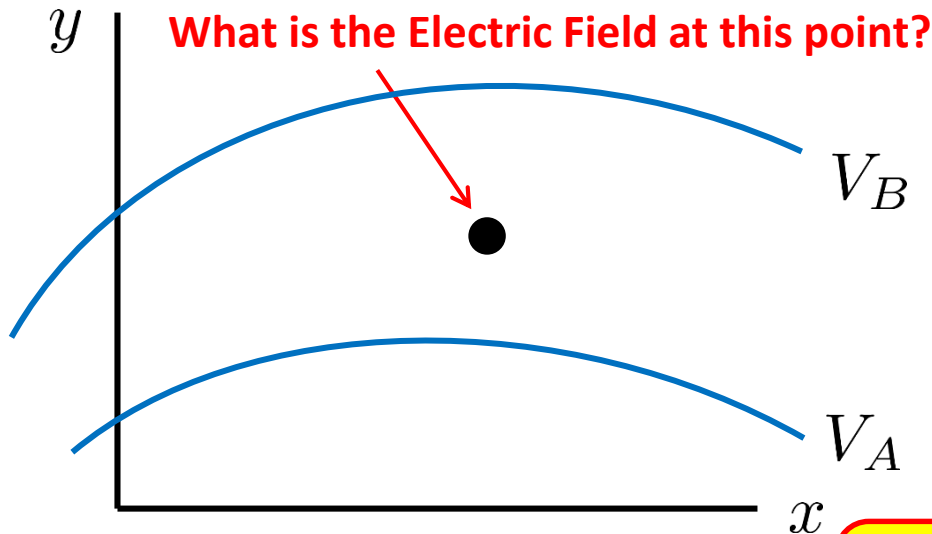
$$\text{So: } \Delta V = V_B - V_A = -37,500 \text{ V}$$

Note: I don't know what  $V = Ed$  gives here!

# Going from the Potential to the Electric Field

*i.e.* If we know  $V$  in a region of space, how do we find  $\vec{E}$  ?

What does it mean to know the potential in a region of space? **Know the equipotentials.**



Recall from PHY181 for a force and potential energy (in 1D):

$$F_x = -\frac{dU}{dx}$$

$$\text{and } \vec{F} = q\vec{E} \quad U = qV$$

$$\text{So: } E_x = -\frac{dV}{dx}$$

**In two dimensions, this becomes:**

**Meaning:**

A changing  $V \Rightarrow$  there is an  $\vec{E}$

$\vec{E}$  points in the direction of decreasing  $V$

$\vec{E}$  is everywhere  $\perp$  to equipotential surfaces

$$E_x = -\frac{\partial V}{\partial x} \quad \& \quad E_y = -\frac{\partial V}{\partial y}$$

**Partial Derivatives?**

# A Quick Lesson in Partial Derivatives

- If you have taken Calc III, then you're familiar with partial derivatives, but if you haven't, here's a quick lesson – *it's one of the easiest things in Calculus*.
- In Calc I, you work with functions of a single variable, e.g.  $f(x)$ , and the derivative of  $f(x)$  is the rate of change of  $f$  as  $x$  changes – called an **ordinary derivative**:

So for  $f(x)$ , the derivative,  $\frac{df(x)}{dx}$  = rate of change of  $f$  as  $x$  changes

- In Calc III, you'll work with functions of more than one variable, e.g.  $g(x,y)$ . How do we handle derivatives for a function like this – **partial derivatives**.

So for  $g(x,y)$ , the partial derivative,  $\frac{\partial g(x,y)}{\partial x}$  = rate of change of  $g$  as  $x$  changes

and, the partial derivative,  $\frac{\partial g(x,y)}{\partial y}$  = rate of change of  $g$  as  $y$  changes

- Note, the notation: a regular “ $d$ ” for ordinary derivative, a “*curly-cue d*” for the partial derivative.
- **Now, when taking partial derivative with respect to  $x$ , what do you do with  $y$ ?  
Treat it as a constant! And for the partial with respect to  $y$ , you treat  $x$  as a constant.**
- Here's a quick example: suppose

$$g(x,y) = 6x^2y + 2xy^3$$

$$\frac{\partial g}{\partial x} = 12xy + 2y^3 \quad \text{and} \quad \frac{\partial g}{\partial y} = 6x^2 + 6xy^2$$

## Whiteboard Problem: 26-2

The electric potential in a region of space is

$$V(x,y) = (150x^2 - 200y^2) \text{ Volts}$$

where  $x$  and  $y$  are in meters.

**What is the electric field in component form at  $(x,y) = (2.0\text{m}, 2.0\text{m})$ ? (LC)**

Remember: Instructions for entering a vector in component form on LC:

For the vector:

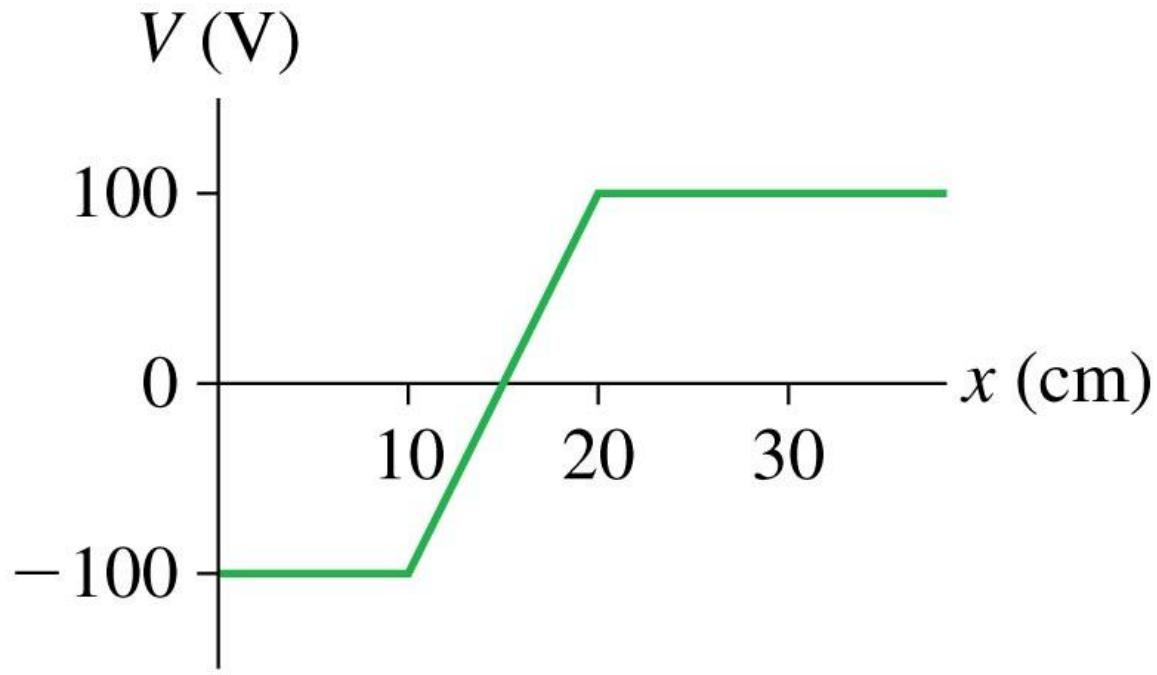
$$\vec{E} = -5.32 \times 10^4 \hat{i} + 1.78 \times 10^3 \hat{j}$$

**Use only two significant figures and the pull down menu for powers of ten if necessary; put the components in parentheses, and use just  $i$  and  $j$  for unit vectors – no hats!**

For the vector above, enter:  **$(-5.3 \cdot 10^4)i + (1.8 \cdot 10^3)j$**

## Whiteboard Problem: 26-3

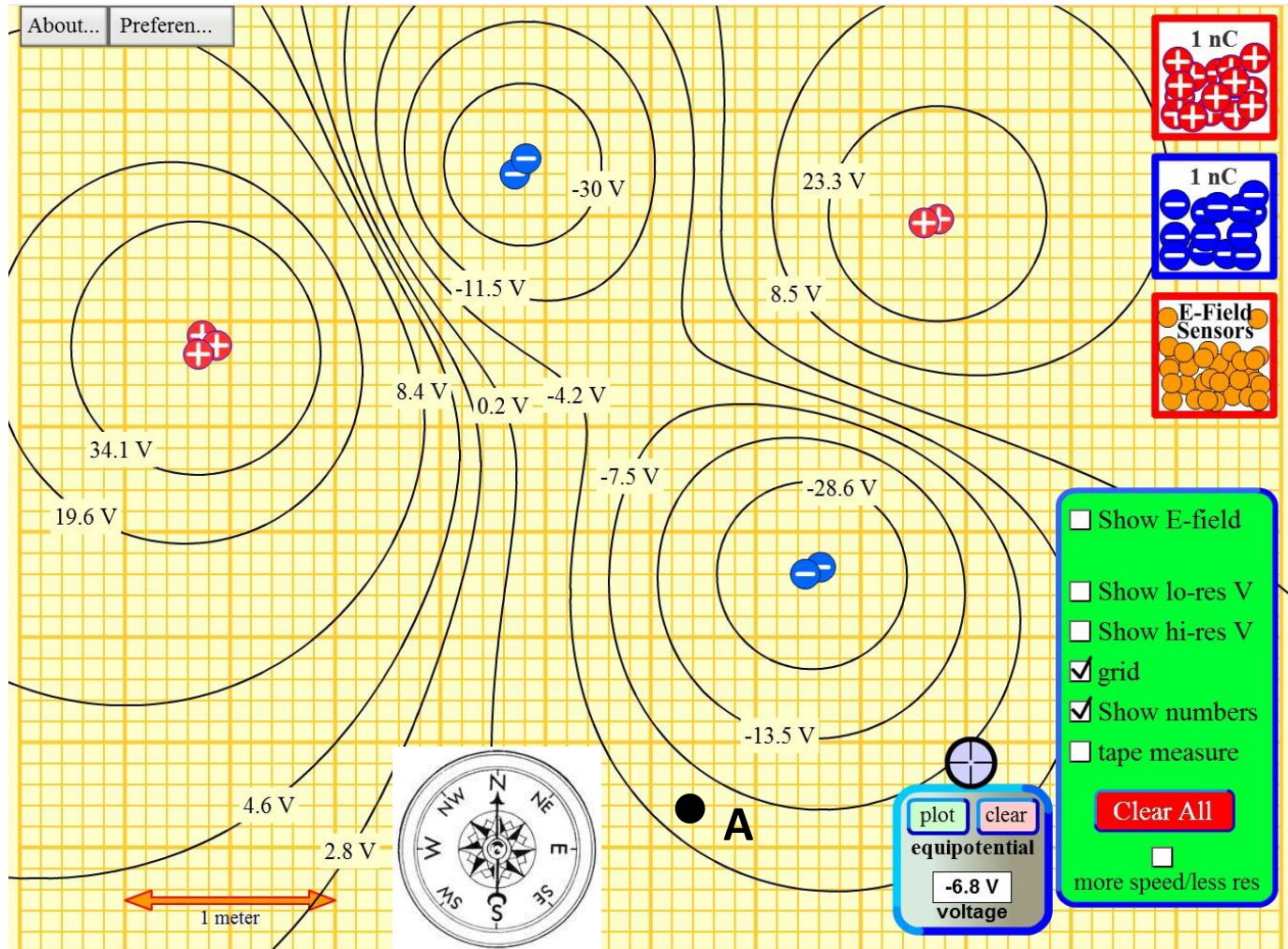
The figure below shows a graph of  $V$  versus  $x$ .  
**Draw the corresponding graph of  $E_x$  versus  $x$ . (LC)**



# Whiteboard Problem: 26-4

The figure shows the equipotentials for a group of charges.

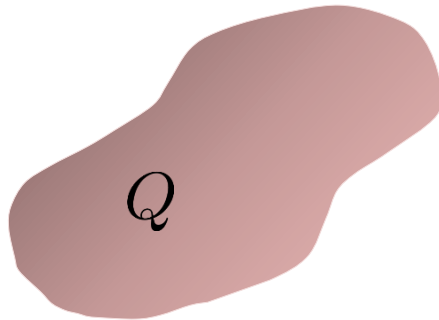
Find an estimate for the strength of the electric field and its direction at the point labelled A. (LC) (give direction as compass direction – see the compass?)



# Using the Potential to find the Electric Field

As we saw in Chapter 25, in most cases, it's easier to find the potential for a given distribution of charge. If that is the case, **then it may be easier to get the field from the potential, instead of integrating directly for the field.**

e.g.



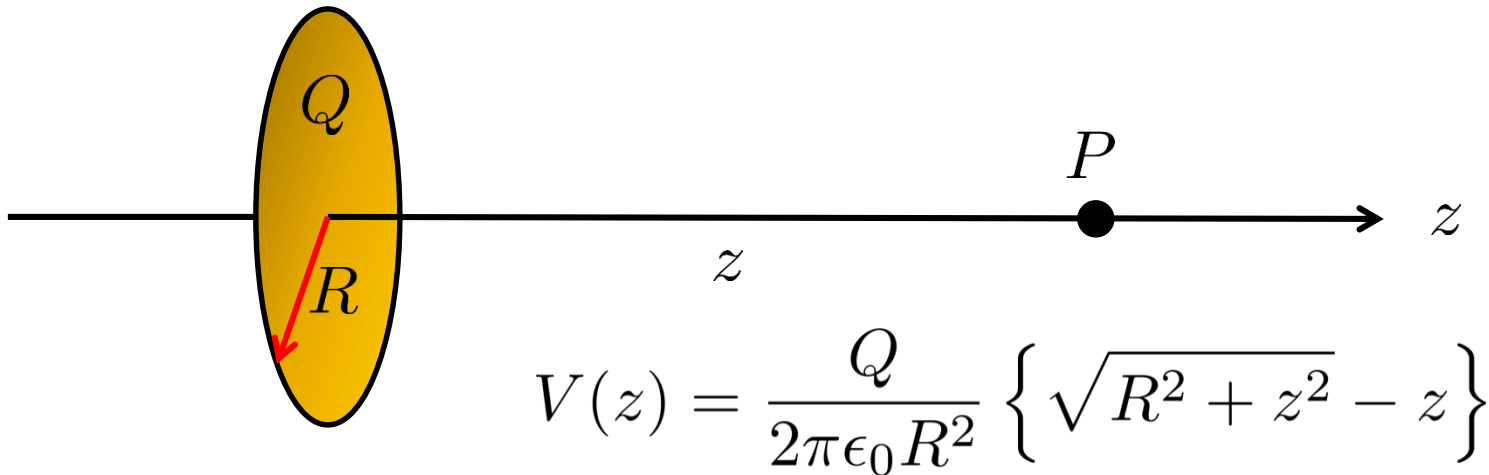
**Suppose we know  $V(x,y)$  in the region around the point P.**

Then, we can get the electric field at P by:

$$E_x = - \left. \frac{\partial V}{\partial x} \right|_{\text{at } P} \quad \& \quad E_y = - \left. \frac{\partial V}{\partial y} \right|_{\text{at } P}$$

## Whiteboard Problem: 26-5

In your text, it was found in Example 25-11 that the potential on the axis of a charged disk is:

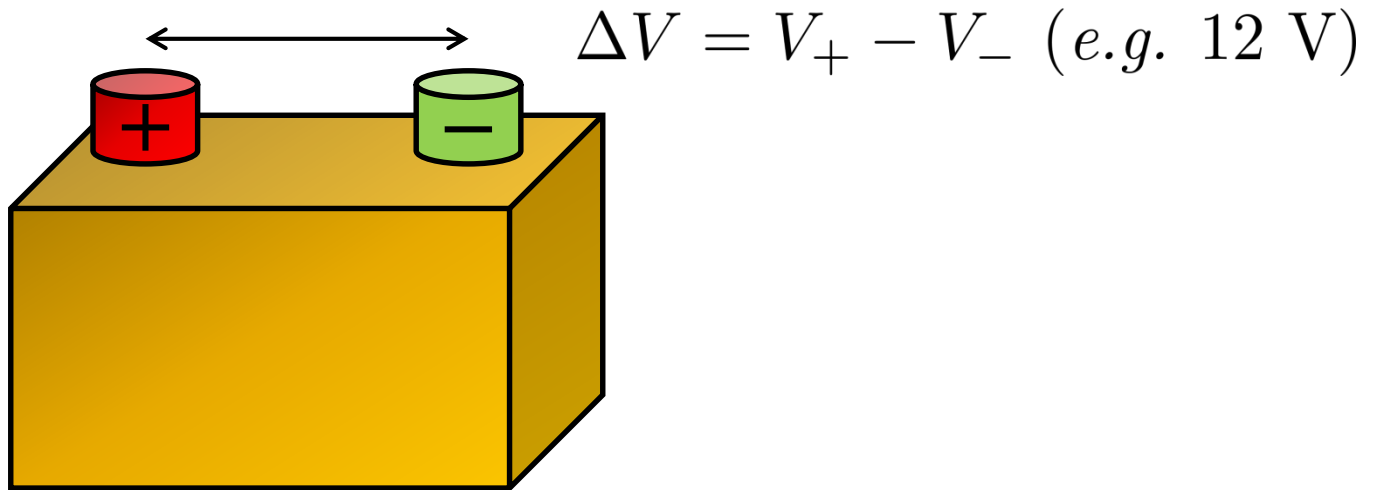


Use this expression for the potential to find an expression for the  $z$ -component of the electric field on the axis. (LC)

# Sources of Potential (EMF)

A source of potential difference is sometimes called *a source of EMF*, a widely used term, which stands for **ElectroMotive Force**. Your author points out that this is an outdated term since Force may or may not be involved when there is a potential difference – but you still see it used frequently.

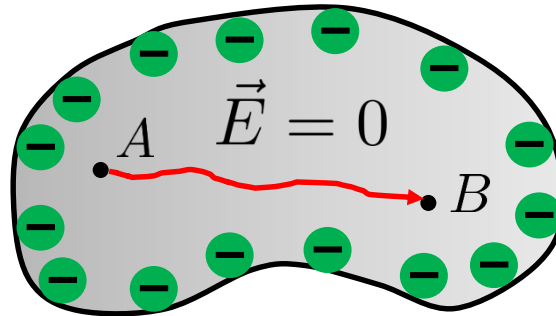
Most sources of potential (or emf) create a potential difference by separating charge either mechanically (e.g. rubbing a rod or a Van de Graff generator) or chemically as in a **battery**:



A lot of interesting chemistry goes on inside a battery, but for our purposes, we can just consider it to be a **source of constant potential difference between the battery terminals**. ([here's an optional video that details the chemistry in a zinc copper battery: https://www.youtube.com/watch?v=7b34XYgADIM&t=791s](https://www.youtube.com/watch?v=7b34XYgADIM&t=791s))

# Revisit: The Charged Conductor in Equilibrium

In the Chapter 24 Readings, we saw that a conductor in equilibrium has **zero electric field** inside, and any excess charge resides on the exterior surface:

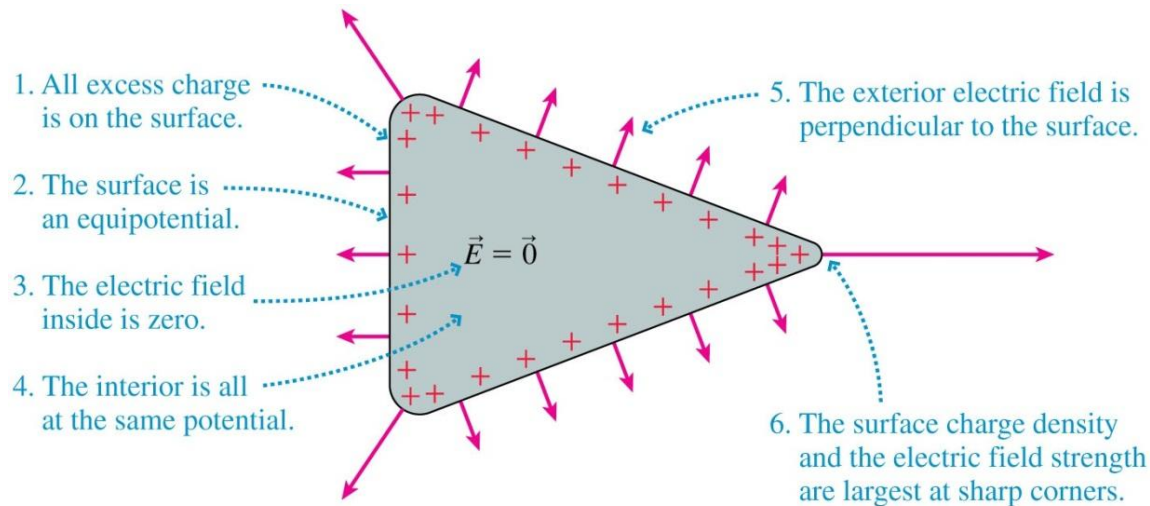


For any two points inside the conductor:

$$\Delta V = V_B - V_A = - \int_A^B \vec{E} \cdot d\vec{s} = 0 \text{ since } \vec{E} = 0 \text{ at all points on any path}$$

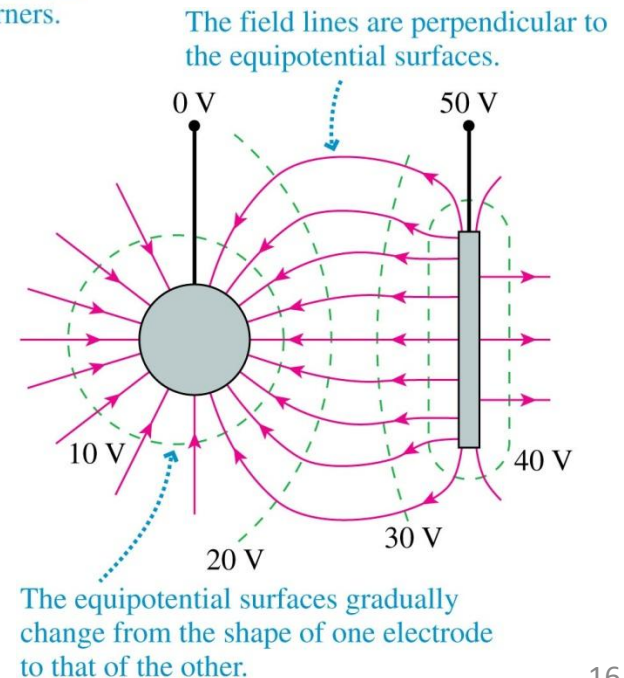
**Therefore, all points inside a conductor in equilibrium are at the same potential, and the surface is an equipotential surface.**

# Summary: The Charged Conductor in Equilibrium



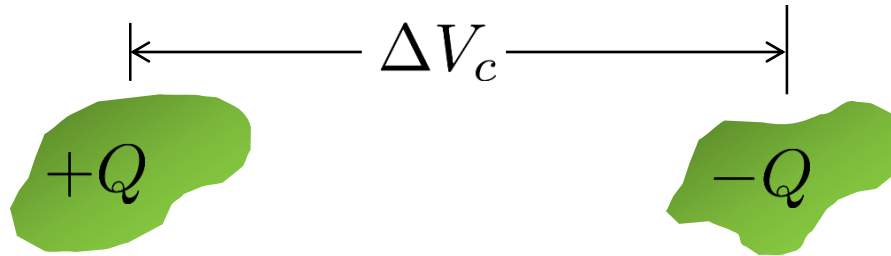
**Important facts for a charged conductor in equilibrium**

**The electric field and equipotentials between two charged conductors**



# Capacitors and Capacitance

We have already introduced the parallel plate capacitor as a source of uniform electric field, **but in general, a capacitor is any two conductors (electrodes) carrying equal but opposite charges:**



There are complicated electric fields and equipotentials between the conductors that depend on their geometry; however, the potential difference is proportional to the charge:

$$\Delta V_c \propto Q \Rightarrow \Delta V_c = \left( \frac{1}{C} \right) Q$$

So, we define the **Capacitance**  
**of the two electrodes as:**

$$C = \frac{Q}{\Delta V_c} \quad \text{Units: } = \frac{C}{V} \equiv \text{Farad (F)}$$

The **capacitance of a capacitor** is the proportionality constant between the charge on the conductors and the potential between them. It depends on the geometry of the conductors.

# PhET Capacitor Demo

Remember from watching the video, **The Story of Electricity Part 1**, that the **Leyden Jar** was a device to store charge - - maybe you've built your own **Leyden jar**. The modern equivalent of the Leyden Jar is the Capacitor. In this sense, the Capacitance of any Capacitor is a measure of how much charge the device can store for a given potential difference, *i.e. it's charge Capacity*.

**In this demo, we will explore how the characteristics of a parallel plate capacitor affect how much charge it can store, i.e. it's capacitance.**

After we explore a bit, we'll develop the necessary physical relations.

The PhET simulation is at:

<https://phet.colorado.edu/en/simulation/capacitor-lab-basics>

***By the way: we'll cover multiple capacitors in combination in chapter 28 (circuits) when it is used.***

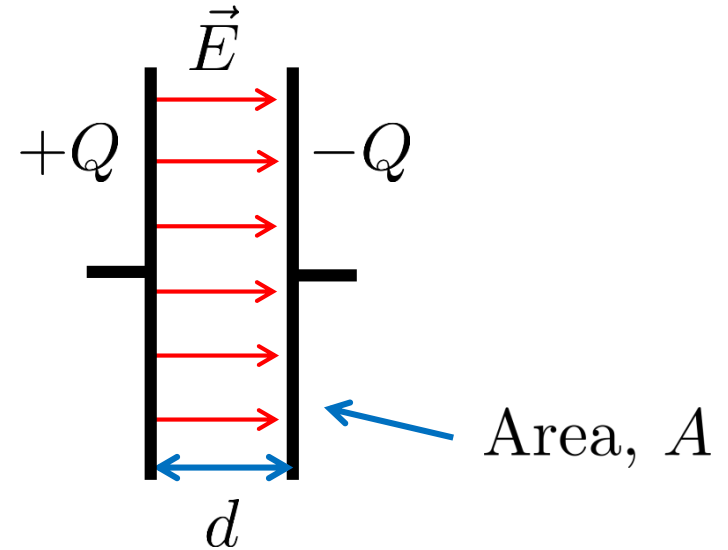
# The Capacitance of a Parallel Plate Capacitor

In the PhET demo, we found that for a **parallel plate capacitor**, the capacitance is larger for larger plate area and smaller plate separation.

**Here, we'll analyze this capacitor geometry:**

For a parallel plate capacitor:

Know:  $E = \frac{Q}{\epsilon_0 A}$  (Uniform)



And for a uniform electric field:

$$\Delta V = Ed = \left( \frac{Q}{\epsilon_0 A} \right) d \quad \text{So:} \quad Q = \left( \frac{\epsilon_0 A}{d} \right) \Delta V$$

For a Parallel Plate Capacitor:  $C = \frac{Q}{\Delta V} = \frac{\epsilon_0 A}{d}$

*As you saw in the PhET demo, the plates can hold more charge for larger  $A$  and/or smaller  $d$ .*

## Whiteboard Problem: 26-6 **If Time**

You need to construct a 100 pF capacitor for a science project. You plan to cut two  $L \times L$  metal squares and insert small spacers between their corners to make a parallel plate capacitor. The thinnest spacers you have are 0.20 mm thick.

**What value of  $L$  will give you the capacitance you want? (LC)**

# Energy Stored in a Capacitor

Your author shows that it takes energy to charge a capacitor, and this energy is stored in the capacitor:

Energy Stored in  
a Capacitor

$$U_c = \frac{Q^2}{2C} = \frac{1}{2}C (\Delta V_c)^2$$

**Where, exactly, does this energy reside?** Consider a parallel plate capacitor:

$$U_c = \frac{1}{2}C (\Delta V_c)^2 = \frac{1}{2} \left( \frac{\epsilon_0 A}{d} \right) (Ed)^2 = \frac{1}{2} \epsilon_0 Ad E^2$$

Now, consider the **Energy Density** - *i.e. the energy per unit volume*:

$$u_E = \frac{U_c}{\text{Volume}} = \frac{\frac{1}{2} \epsilon_0 Ad E^2}{Ad} = \frac{1}{2} \epsilon_0 E^2$$

So, the energy is stored in the field with energy density:

$$u_E = \frac{1}{2} \epsilon_0 E^2$$

***This is a general result that is true for all electric fields.***